

PHY V2500: QUANTUM MECHANICS I

Problem Set 1

Due September 11, 2025

Problem 1

In class I calculated the elements $X_{nn'}$. In a similar way, we define

$$P_{nn'} = \int_0^L dx v_n(x) \left(-i\hbar \frac{\partial}{\partial x} \right) v_{n'}(x)$$

where

$$v_n(x) = \sqrt{\frac{2}{L}} \sin(n\pi x/L), \quad x \in [0, L]$$

Carry out the integration to obtain $P_{nn'}$. Also show that $P_{nn'}^* = P_{n'n}$.

Solution

Carrying out the indicated differentiation,

$$\begin{aligned} P_{nn'} &= (-i\hbar) \frac{2}{L} \frac{n'\pi}{L} \int_0^L dx \sin(n\pi x/L) \cos(n'\pi x/L) \\ &= (-i\hbar) \frac{n'\pi}{L^2} \int_0^L [\sin((n+n')\pi x/L) + \sin((n-n')\pi x/L)] \\ &= (-i\hbar) \frac{n'\pi}{L^2} \frac{L}{\pi} \left[-\frac{\cos((n+n')\pi x/L)}{n+n'} - \frac{\cos((n-n')\pi x/L)}{n-n'} \right]_0^L \\ &= -i \frac{2\hbar}{L} \frac{nn'}{n^2 - n'^2} [1 - (-1)^{n+n'}] \end{aligned}$$

This result applies for $n \neq n'$. For $n = n'$, we have

$$P_{nn} = (-i\hbar) \frac{n'\pi}{L^2} \int_0^L \sin(2\pi n x/L) = (-i\hbar) \frac{n'\pi}{L^2} \frac{L}{2\pi n} [-\cos(2\pi n x/L)]_0^L = 0$$

Taking the complex conjugate of $P_{nn'}$, we see

$$P_{nn'}^* = i \frac{2\hbar}{L} \frac{nn'}{n^2 - n'^2} [1 - (-1)^{n+n'}] = -i \frac{2\hbar}{L} \frac{nn'}{n'^2 - n^2} [1 - (-1)^{n+n'}] = P_{n'n}$$

Problem 2

Consider the function

$$\psi(x) = C x \sin(\pi x/L)$$

which is defined on the interval $[0, L]$.

a) Identify the value of C such that ψ is normalized, i.e., such that

$$\int_0^L dx \psi^* \psi = 1$$

b) The function ψ is square-integrable over the interval $[0, L]$ and vanishes at $x = 0, L$. So it can be expanded in terms of the basis $\{v_n(x)\}$ given in Problem 1 as

$$\psi(x) = \sum_n c_n v_n(x)$$

Calculate the coefficients c_n .

Solution

a) \mathcal{C} is determined by

$$1 = |\mathcal{C}|^2 \int_0^L dx x^2 \sin^2(\pi x/L) = \frac{|\mathcal{C}|^2}{2} \int_0^L dx x^2 (1 - \cos(2\pi x/L))$$

We note that

$$\frac{\partial}{\partial x} \left[\frac{x^3}{3} - \left(\frac{x^2 \sin \alpha x}{\alpha} + \frac{2x \cos \alpha x}{\alpha^2} - \frac{2 \sin \alpha x}{\alpha^3} \right) \right] = x^2 - x^2 \cos \alpha x$$

Taking $\alpha = 2\pi/L$, we find

$$\begin{aligned} 1 &= \frac{|\mathcal{C}|^2}{2} \left[\frac{x^3}{3} - \left(\frac{x^2 \sin \alpha x}{\alpha} + \frac{2x \cos \alpha x}{\alpha^2} - \frac{2 \sin \alpha x}{\alpha^3} \right) \right]_0^L \\ &= \frac{|\mathcal{C}|^2}{2} L^3 \left(\frac{1}{3} - \frac{1}{2\pi^2} \right) \end{aligned}$$

This identifies \mathcal{C} as

$$\mathcal{C} = \frac{1}{L^{\frac{3}{2}}} \left(\frac{1}{6} - \frac{1}{4\pi^2} \right)^{-\frac{1}{2}}$$

b) From the orthonormality of v_n , we have

$$c_n = \int_0^L dx v_n(x) \psi(x) = \sqrt{\frac{2}{L}} \mathcal{C} \int_0^L dx x \sin(n\pi x/L) \sin(\pi x/L)$$

We can now write

$$\sin(n\pi x/L) \sin(\pi x/L) = \frac{1}{2} [\cos((n-1)\pi x/L) - \cos((n+1)\pi x/L)]$$

We then note that

$$\frac{\partial}{\partial x} \left[\frac{x \sin \alpha x}{\alpha} + \frac{\cos \alpha x}{\alpha^2} \right] = x \cos \alpha x$$

Using these two results, we find

$$\begin{aligned} c_n &= \frac{\mathcal{C}}{\sqrt{2L}} \left[\frac{L x \sin((n-1)\pi x/L)}{\pi (n-1)} + \frac{L^2 \cos((n-1)\pi x/L)}{\pi^2 (n-1)^2} \right. \\ &\quad \left. - \frac{L x \sin((n+1)\pi x/L)}{\pi (n+1)} + \frac{L^2 \cos((n+1)\pi x/L)}{\pi^2 (n+1)^2} \right]_0^L \\ &= \frac{\mathcal{C} L^2}{\sqrt{2L}} (1 - (-1)^{n+1}) \left[\frac{1}{(n+1)^2} - \frac{1}{(n-1)^2} \right] \end{aligned}$$

This applies for $n \neq 1$. For $n = 1$, we have, with $\alpha = 2\pi/L$,

$$\begin{aligned} c_1 &= \sqrt{\frac{2}{L}} \frac{C}{2} \int dx x (1 - \cos(2\pi x/L)) = \sqrt{\frac{2}{L}} \frac{C}{2} \left[\frac{x^2}{2} - \left(\frac{x \sin \alpha x}{\alpha} + \frac{\cos \alpha x}{\alpha^2} \right) \right]_0^L \\ &= C \sqrt{\frac{1}{8L}} \end{aligned}$$

Problem 3

Consider a Hilbert space and a basis of vectors $\{|e_i\rangle\}$, $i = 1, 2, 3, \dots$. Unlike what we did in class, we will not assume that these are orthonormal. Thus

$$\langle e_i | e_j \rangle = g_{ij}$$

where g_{ij} is not δ_{ij} , but can be taken to be invertible.

- Show that the matrix \mathbf{g} (with matrix elements g_{ij}) is hermitian.
- Expand an arbitrary vector in terms of this basis as

$$|a\rangle = \sum_i a_i |e_i\rangle$$

Calculate the coefficients a_i in terms of the inner product $\langle e_i | a \rangle$.

- Using the solution for part b), obtain the completeness relation for the basis $\{|e_i\rangle\}$.

Solution

- From the properties of the Hilbert space discussed in class,

$$\langle a | b \rangle = \langle b | a \rangle^*$$

Thus

$$g_{ij} = \langle e_i | e_j \rangle = \langle e_j | e_i \rangle^* = g_{ji}^* = (g^T)^*_{ij} = (g^\dagger)_{ij}$$

Thus we have $\mathbf{g}^\dagger = \mathbf{g}$, showing that \mathbf{g} is hermitian.

- Taking the inner product of $|a\rangle$ with $\langle e_j |$, we get

$$\langle e_j | a \rangle = \sum_i a_i \langle e_j | e_i \rangle = \sum_i g_{ji} a_i$$

Multiplying by $(g^{-1})_{kj}$ and summing over j , we find

$$\sum_j (g^{-1})_{kj} \langle e_j | a \rangle = \sum_{j,i} (g^{-1})_{kj} g_{ji} a_i = \sum_i \delta_{ki} a_i = a_k$$

Thus $a_k = \sum_j (g^{-1})_{kj} \langle e_j | a \rangle$, or $a_i = \sum_j (g^{-1})_{ij} \langle e_j | a \rangle$. Using this in the expansion for $|a\rangle$,

$$|a\rangle = \sum_i |e_i\rangle \sum_j (g^{-1})_{ij} \langle e_j | a \rangle = \left(\sum_{i,j} |e_i\rangle (g^{-1})_{ij} \langle e_j | \right) |a\rangle$$

Since this applies for any $|a\rangle$, we get the relation

$$\sum_{i,j} |e_i\rangle (g^{-1})_{ij} \langle e_j| = \mathbb{1}$$

This is the completeness relation when we do not have orthonormality.

Problem 4

We have considered the basis $\{v_n\}$ as in Problem 1. Now consider a different set of functions defined on the interval $[0, L]$ as

$$u_n(x) = \sqrt{\frac{2}{L}} \cos(n\pi x/L), \quad n = 1, 2, \dots$$

These functions do not vanish at the boundaries. Rather they obey $u'_n(x) = 0$ at $x = 0, L$. (These are what are called Neumann boundary conditions.) Carry out the needed integration to identify the integral

$$\int_0^L dx u_n(x) u_{n'}(x)$$

Solution

We write the integral as

$$\begin{aligned} \int_0^L dx u_n(x) u_{n'}(x) &= \frac{2}{L} \int_0^L dx \cos(n\pi x/L) \cos(n'\pi x/L) \\ &= \frac{1}{L} \int_0^L dx \left[\cos[(n+n')\pi x/L] + \cos[(n-n')\pi x/L] \right] \\ &= \frac{1}{L} \frac{1}{\pi} \left[\frac{\sin[(n+n')\pi x/L]}{n+n'} + \frac{\sin[(n-n')\pi x/L]}{n-n'} \right]_0^L = 0 \end{aligned}$$

This applies for $n \neq n'$. For $n = n'$, we get

$$\int_0^L dx u_n(x) u_{n'}(x) = \frac{1}{L} \int_0^L dx \left[\cos[2\pi n x/L] + 1 \right] = 1$$

Combining these two results we have

$$\int_0^L dx u_n(x) u_{n'}(x) = \delta_{nn'}$$
